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1 **Effects of frozen storage temperature on the elasticity of tendons from a small**  
2 **murine model**

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12 Running head: Effects of frozen storage temperature on tendons

13 Implications: Ultra-low storage temperature alters (1) the long-range order of collagen  
14 packing in fibrils, resulting in fibril rupture at higher stresses, (2) the composition of  
15 extra-fibrillar matrix, resulting in an increase in the interaction energy between fibrils via  
16 collagen-bound proteoglycans. Consequently, the tissue strength, stiffness, and fracture  
17 toughness increase.

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20

21 **Abstract**

22 The basic mechanism of reinforcement in tendons addresses the transfer of stress,  
23 generated by the deforming proteoglycan-rich (PG) matrix, to the collagen fibrils.  
24 Regulating this mechanism involves the interactions of PGs on the fibril with those in the  
25 surrounding matrix and between PGs on adjacent fibrils. This understanding is key to  
26 establishing new insights on the biomechanics of tendon in various research domains.  
27 However, the experimental designs in many studies often involved long sample  
28 preparation time. To minimise biological degradation the tendons are usually stored by  
29 freezing. Here, we have investigated the effects of commonly used frozen storage  
30 temperatures on the mechanical properties of tendons from the tail of a murine model  
31 (C57BL6 mouse). Fresh (unfrozen) and thawed samples, frozen at temperatures of  $-20\text{ }^{\circ}\text{C}$   
32 and  $-80\text{ }^{\circ}\text{C}$ , respectively, were stretched to rupture. Freezing at  $-20\text{ }^{\circ}\text{C}$  revealed no effect  
33 on the maximum stress ( $\sigma$ ), stiffness ( $E$ ), the corresponding strain ( $\epsilon$ ) at  $\sigma$ , and strain  
34 energy densities up to  $\epsilon$  ( $u$ ) and from  $\epsilon$  until complete rupture ( $u_p$ ). On the other hand,  
35 freezing at  $-80\text{ }^{\circ}\text{C}$  led to higher  $\sigma$ ,  $E$  and  $u$ ;  $\epsilon$  and  $u_p$  were unaffected. The results  
36 implicate changes in the long-range order of radially-packed collagen molecules in  
37 fibrils, resulting in fibril rupture at higher stresses, and changes to the composition of  
38 extra-fibrillar matrix, resulting in an increase in the interaction energy between fibrils via  
39 collagen-bound PGs.

40

## 41 **Introduction**

42 This paper is concerned with the effects of frozen storage temperature on the elasticity of  
43 tendons, which are biological examples of fibre reinforced composites comprising highly  
44 paralleled collagen fibrils embedded in the hydrated proteoglycan-rich (**PG**) extracellular  
45 matrix (**ECM**) (Hukins & Aspden 1985; Figure 1A). The basic mechanism of  
46 reinforcement in tendons addresses the transfer of stress generated in the deforming PG-  
47 rich matrix, in response to an external load, to the collagen fibrils (Goh *et al.*, 2005).  
48 Regulating this mechanism involves the interactions of PGs on the fibrils with those in  
49 the surrounding matrix and between adjacent fibrils (Goh *et al.*, 2007). The stress transfer  
50 mechanism is fundamental to understanding the biomechanics of load transmission, from  
51 muscle to bone by tendons, in many studies. These studies may be found in biomedical  
52 science, e.g. addressing the biomarkers of ageing (Derwin & Soslowsky, 1999), injury  
53 and repair (Lin *et al.*, 2004), and meat science, e.g. investigating the mechanical response  
54 in relation to the quality of the attachment of muscle to bone from livestock (Moussa *et al.*  
55 *et al.*, 2007; Gondret *et al.*, 2009). Unfortunately, experimental designs in these studies  
56 often involve long sample preparation time. To minimise biological degradation tendons  
57 are stored by freezing, usually at around -20 °C (Moussa *et al.*, 2007; Goh *et al.*, 2008;  
58 Gondret *et al.*, 2009) or at ultra-low temperatures, e.g. -80 °C (Giannini *et al.*, 2008),  
59 until the day of testing. Although the processes that regulate the formation of ice  
60 crystallites and the disruption of cells and ECM organisation are well-known (Muldrew  
61 & McGann, 1990), how frozen storage temperature affects the basic mechanisms that  
62 regulate the ECM elasticity, leading to changes in the mechanical properties, is not clear.

63 We have investigated the influence of frozen storage temperatures on the mechanical  
64 properties of tendons from the tail of a small murine model (C57BL6 mouse). Here we  
65 are concerned with the maximum stress,  $\sigma$ , on the stress-strain curve, stiffness ( $E$ ), strain  
66 ( $\epsilon$ ) at  $\sigma$ , strain energy densities up to  $\epsilon$  ( $u$ ) and from  $\epsilon$  until complete rupture ( $u_p$ ) (Goh *et*  
67 *al.*, 2005). Since the genome of the C57BL6 mouse has been sequenced, it has been  
68 regarded as an ideal murine model in the context of integrative biology study; it is widely  
69 used for laboratory experiments to illuminate basic mechanisms which could address  
70 similar outcomes in other mammals (Laissue *et al.*, 2009). In particular, tissues from the  
71 murine model have been used extensively to evaluate biomechanical properties to study  
72 systemic changes caused by injury and during healing (Lin *et al.*, 2004), ageing of ECM  
73 (Goh *et al.*, 2008) and disruption of the signal pathway for regulating ECM in genetically  
74 engineered mice that produced alterations in collagen (Derwin & Soslowsky, 1999) and  
75 PGs (Derwin & Soslowsky, 1999). Application of the principles of fibre reinforced  
76 composites to connective tissues such as tendons has yielded insights on the mechanisms  
77 of elastic and plastic stress transfer in collagen fibrils and on fibril-fibril interactions  
78 across the hydrated PG-rich matrix within ECM for regulating tissue elasticity (Goh *et*  
79 *al.*, 2005). Here, these insights will be used to speculate on the sensitivity of these  
80 mechanisms to low storage temperatures.

## 81 **Material and Methods**

### 82 *Sample preparation*

83 Tendons, of about 200 microns thick, were dissected from the tail of a 16 month-old  
84 mouse (C57BL6) which was sacrificed in accordance with the procedure laid down by

85 the National Research Council for the care and use of laboratory animals. The tendons  
86 were transected and divided equally into three groups consisting of ten samples per  
87 group. One group was designated as control and was tested immediately. Samples from  
88 the second (**M20**) and third (**M80**) groups were immediately placed in freezers operating  
89 at -20 °C (MDF-U5411, Sanyo) and -80 °C (MDF-U73V, Sanyo), respectively. After 24  
90 hours the frozen samples were thawed at room temperature for analysis.

### 91 *Mechanical tests*

92 A small-scale horizontal tensile rig, designed and built in-house, was used for mechanical  
93 tests (Goh *et al.*, 2008). Each sample was mounted onto the rig, secured by the grips at its  
94 ends and hydrated by submerging in phosphate buffer saline, pH 7.2, at room  
95 temperature. The rig was mounted onto an inverted microscope (DM-IL, Leica) for  
96 sample observation during the test. A displacement rate of 0.067 mm/s was applied to  
97 stretch the sample until the tissue ruptured (Goh *et al.*, 2008). The force generated in the  
98 sample was recorded by a load-cell (UF-1, Pioden Controls Ltd) coupled to one grip; the  
99 sample extension corresponded to the relative displacement of the grips, which was  
100 recorded by a linear variable displacement transducer (CDP-M, Tokyo Sokki Kenkyujo).

### 101 *Determination of mechanical properties*

102 Stress-strain plots were obtained from the load-extension data. Here, stress is defined as  
103 the ratio of the load to the nominal cross-sectional area of the sample; strain is the ratio of  
104 the change in the sample length to the length described by the grip-to-grip distance. By  
105 fitting an appropriate polynomial equation to the data points on the plot, from stress = 0

106 to  $\sigma$ , the point of inflexion was identified (Figure 1B). The gradient at this point  
107 parameterised  $E$  (Goh *et al.*, 2008).

108 From the area under the stress-strain curve (Figure 1B), the strain energy density ( $u$ ) up to  
109  $\sigma$  was evaluated to determine the energy needed to cause tissue failure; in this region the  
110 energy absorbed is dominated by fibril recruitment and fibril yielding (Goh *et al.*, 2008).

111 We have also determined the strain energy density  $u_p$ , from  $\epsilon$  until when the tissue  
112 ruptured completely, to account for the energy absorbed by the dominant failure modes,  
113 namely fibril pull-out by frictional sliding at the fibril-PG matrix interface, as well as PG  
114 matrix rupture (Sikoryn & Hukins, 1988).

#### 115 *Statistical analyses*

116 Representative (mean  $\pm$  standard error (s.e.)) values of the mechanical parameters, i.e.  $\sigma$ ,  
117  $E$ ,  $\epsilon$ ,  $\epsilon_y$ , and  $u$ , for each treatment were determined. Paired sample t-test was used to  
118 analyse for differences between the control versus each treatment group at  $\alpha = 0.05$ .  
119 Differences due to the treatment were considered significant if the P value was less than  
120 0.05.

#### 121 **Results**

122 Figure 2 shows stress-strain plots derived from the control, M20 and M80 groups. In  
123 general, during the loading stage, all samples featured a toe region followed by a rapid  
124 increase in the rate of change of stress with respect to strain; during tissue failure, i.e.  
125 beyond the strain at maximum stress, the gradients decreased rapidly until rupture. Of the

126 three groups, the M80 group yielded the largest gradient during the loading stage and  
127 during tissue failure.

128 Examination of the micro-structure of tendons from the control, M20 and M80 groups  
129 revealed wavy collagen fibres (which are one hierarchy above collagen fibrils) which  
130 exhibited the usual structural changes, i.e. de-crimping and straightening (Figure 3a),  
131 during loading. All the three groups responded with similar structural changes during  
132 mechanical loading. Immediately after  $\sigma$ , the tendon began to fail (Figure 3b). Here, a  
133 thin section bridging the bulk of tissue was observed (Arrow A). Rupture and pull-out of  
134 collagen fibres were also observed; the ends of these structures are indicated by arrow B.  
135 At the lower hierarchical level, this could implicate collagen fibril rupture and pull-out.  
136 Longitudinal splits in the tendon were also observed, implicating the rupture of the PG  
137 matrix (arrow C).

138 The results of  $\sigma$ ,  $E$ ,  $\epsilon$ ,  $u_p$  and  $u$  for the control, M20 and M80 groups are listed in Table 1.  
139 In all the cases considered here, no significant difference was observed between the  
140 control and M20 groups. No significant difference was observed between the control and  
141 M80 groups for  $\epsilon$  and  $u_p$ . However, the M80 group yielded magnitudes of  $\sigma$ ,  $E$  and  $u$ ,  
142 which were 1.4, 1.5 and 1.6 times significantly higher, respectively, than the control  
143 group.

## 144 **Discussion**

145 Collagen fibrils feature a crystalline arrangement of collagen molecules. This crystalline  
146 arrangement is characterised by a long-range order of radially-packed (side-to-side)  
147 collagen molecules and an ordered array of axially staggered collagen molecules along

148 the fibril axis (Laing *et al.*, 2003). The long-range order characterises the crystalline  
149 system in which portions deep within the fibril, i.e. far from the fibril-PG matrix  
150 interface, exhibit similar strain response; strain energy will not concentrate at the  
151 interface but will eventually be absorbed by molecules at greater depths from the  
152 interface. Decorin PGs associated with the collagen fibrils consist of a core protein,  
153 bound to collagen, with a side chain known as glycosaminoglycans (**GAGs**; Figure 1;  
154 Screen *et al.*, 2006). The non-steric interactions between GAGs, from the respective  
155 decorin PGs on adjacent fibrils, is implicated for regulating the stress transfer to collagen  
156 fibrils as the fibrils slide relative to one another when the tissue is loaded (Screen *et al.*,  
157 2006; Goh *et al.*, 2007). According to the elasticity of materials (Sun, 2009), it follows  
158 that  $E$  and  $\sigma$  may be determined by equilibrating the external mechanical load to the  
159 responses of the inter-molecular forces of the tendons. Thus,

$$E \propto U_1 / V , \quad (1)$$

$$\sigma \propto U_2 / V , \quad (2)$$

160 where  $U_1$  and  $U_2$  are the sum of cohesive energy per bond, at equilibrium between the  
161 biomacromolecules, within a volume,  $V$  (Sun, 2009). Although  $E$  and  $\sigma$  have similar  
162 dimensions, the molecular interactions involved in these parameters are not the same.  
163 Equation (1) is valid whether the tendon is undergoing elastic or plastic deformation,  
164 which is determined by the forces of interactions between the GAGs of decorin PGs on  
165 adjacent collagen fibrils. Equation (2) is only valid when the tendon is experiencing  
166 plasticity, which could involve the plastic deformation of fibrils and the breaking of  
167 covalent bonds between collagen molecules in the fibril. Thereafter, the broken bonds

168 could inhibit the motion of collagen molecules around the dislocations, leading to a  
169 strengthened tissue as energy of the remaining bonds increases and the bond length  
170 contracts (Sun, 2009).

171 It follows that any increase in  $E$  is regulated by the Van Der Waals (**VDW**) forces  
172 between GAGs from collagen-bound PGs, at adjacent fibrils. The interactions are  
173 significant for as long as the fibrils are in close proximity to one another. During  
174 freezing, as water withdraws from the water compartment, beginning in the PG matrix  
175 (Bevilacqua *et al.*, 1979), and precipitates into ice crystallites, the ion concentration  
176 increases in the non-water compartment (Muldrew & McGann, 1990), creating a  
177 concentration gradient across the unfrozen zone. As cations (e.g.  $\text{Na}^+$ ) diffuse to the  
178 anionic charge of the GAGs, the increase in the concentration of cations around the  
179 GAGs increases the attractive forces between the GAGs. This effect may be appreciable  
180 at  $-80\text{ }^\circ\text{C}$  but not at  $-20\text{ }^\circ\text{C}$ . Moreover, the changes occurring in the tendons from the M80  
181 group may not be completely reversible during thawing; in this case, a greater GAG-  
182 GAG interaction increases the resistance for fibril-fibril sliding when the tissue is  
183 stretched so that the tissue becomes stiffer than unfrozen ones. To understand how this  
184 could have happened, a molecular model was developed to study GAG-GAG interactions  
185 (following from Redaelli *et al.*, 2003). Mobile cations, namely  $\text{Na}^+$ , were designated in  
186 closed proximity to each group to model a hypertonic PG matrix environment; the  
187 number of  $\text{Na}^+$  ( $N$ ) was varied to investigate the effects on the VDW energy which was  
188 evaluated as a function of the distance between the PG molecules,  $\delta$ . The model was  
189 solved by molecular mechanics using software Hyperchem (Version 7.5, Hypercube,  
190 Inc.). It was found that VDW energy, which decreased as  $\delta$  increased, was affected by  $N$ .

191 In particular, VDW energy increased by about 1.2 times as  $N$  increased two-fold but the  
192 profile of the energy- $\delta$  curves was not affected. Further details can be found in  
193 supplementary material. Thus we concluded that the overall effect of freezing led to a  
194 localisation and densification of charge and energy. Hence  $E$  has to increase as predicted  
195 by Equation (1).

196 Any increase in  $\sigma$  could be attributed to changes in the packing of collagen molecules in  
197 the fibril. Low storage temperature could alter the long-range order of radial packing of  
198 collagen molecules, which could be assessed by the diffuse scatter region from x-ray  
199 diffraction patterns of these tissues. Changes to the diffuse region occur during freezing;  
200 tissues frozen at  $-150\text{ }^{\circ}\text{C}$  exhibit reduced diffuse regions suggesting that the long range  
201 order is increased (Laing *et al.*, 2003). Whether the effect is reversible upon thawing is  
202 not clear; Hickey and Hukins (1979) had shown that thawed tissues (frozen at  $-35\text{ }^{\circ}\text{C}$ )  
203 exhibited smaller diffuse regions as compared to control samples. An increase in the  
204 long-range order implies that the fine structure, i.e. the crystalline arrangement of  
205 collagen molecules, of the fibril would have to be disturbed to a very much greater depth.  
206 More molecules would be involved in absorbing the energy needed to cause fibril  
207 rupture; thus more energy is absorbed by the tissue. As pointed out earlier in this section,  
208 stresses will not concentrate at regions near to the fibril surface, an important factor  
209 determining the failure of fibrils. Instead, stresses will be distributed deeper within the  
210 fibril, thus making the fibril less susceptible to fracture so that a higher  $u$ , and  
211 consequently higher  $\sigma$  as predicted by Equation (2), would be needed to rupture the  
212 tendon.

213 We note that  $u_p$  is not affected by low storage temperatures. This suggests that low  
214 storage temperature did not have a significant effect on the molecular interaction at the  
215 fibril-PG interface and in the bulk of the PG matrix. We also note that the extensibility of  
216 the tendon, which is parameterised by  $\epsilon$ , was not sensitive to low storage temperatures.  
217 We recalled that x-ray diffraction patterns of stretched tendons revealed that the tissue  
218 strain would always be larger than the strain within the collagen fibrils (Puxkandl *et al.*,  
219 2002). It is possible that the changes affecting the strain in individual fibrils would have  
220 negligible influence on the overall strain at maximum stress in the tissue stored at sub-  
221 zero temperature.

222 The purpose of this paragraph is to discuss the effects of cooling on the size of the ice  
223 crystallites formed in the tendons. According to Bevilacqua *et al.* (1979), a combination  
224 of the effects of solute concentration (due to an increase in viscosity) and thermal  
225 gradients (due to latent heat release and heat lost from the unfrozen zone), in the region  
226 between the crystallites, could lead to super-saturation and super-cooling. Theoretically,  
227 the cooling rate of the tissue is proportional to the square of the moving speed of the  
228 freezing front,  $u^2$ , while the temperature gradient at the solid-liquid interface ( $G$ ) is  
229 proportional to  $u$  (Miyawaki *et al.*, 1992). Here,  $G = [T_m - T_s]/x$  where  $x$  is the length of  
230 the frozen crystallite and  $T_s$  and  $T_m$  are the operating temperature of the freezer and the  
231 temperature at the surface of the crystallite, respectively. Empirically, the thickness ( $d_p$ )  
232 of ice crystallites is related to  $u$  and  $G$ ; this relationship can be represented as

$$d_p \propto 1/\sqrt{uG} \quad (3)$$

233 (Tiller & Rutter, 1956). It follows that the cooling rate is related to  $d_p$  in Equation (3).  
234 Thus, slow freezing rates cause ice to grow in extracellular locations, resulting in large  
235 crystallites and maximum dislocation of water; rapid cooling produces small ice  
236 crystallites in both cells and extracellular matrix (Bevilacqua *et al.*, 1979). In particular,  
237 slow cooling rates correspond to wider spread of sizes; fast cooling rates correspond to  
238 narrower spread of sizes and the crystallite growth is more organised (Pardo *et al.*, 2002).  
239 Consider a model of the crystallite in the tendon. Based on the results of Bevilacqua *et al.*  
240 (1979) and Chevalier *et al.* (2000), using an order of magnitude estimate, we find  $x \sim 100$   
241  $\mu\text{m}$ . Owing to convection in the freezers, we may approximate  $T_s \sim T_f = 0^\circ\text{C}$ , the  
242 freezing temperature of water. With these values, Eq. (3) was evaluated for the upper  
243 limit (0.41 mm/s; Miwayaki *et al.*, 1992) and lower limit (0.06 mm/s; Chevalier *et al.*,  
244 2000) of  $u$ , i.e. corresponding to four combinations, for each case of  $T_m = -20^\circ\text{C}$  and  $-80$   
245  $^\circ\text{C}$ . By determining the ratio  $d_p|_{-20^\circ\text{C}} / d_p|_{-80^\circ\text{C}}$ , this simple analysis predicts that the ice  
246 crystallites formed at the higher freezer temperature are bigger than those formed at the  
247 lower freezer temperature by 1.16 to 3.46 times. In this case, the more organised smaller  
248 crystallites are dispersed uniformly within the inter-fibrillar matrix, raising the ion  
249 concentration around the GAGs. This argument complements an earlier suggestion that  
250 the increase in GAG interactions, arising from a localisation and densification of charge  
251 and energy around the GAGs associated with collagen fibrils, yields a higher  $E$ .  
252 In conclusion, ultra-low storage temperature affects only  $\sigma$ ,  $E$  and  $u$  in tail tendons from  
253 the C57BL6 mouse, leading to tissues with larger  $\sigma$ ,  $E$  and  $u$ . The results of this study  
254 imply that care is required when considering the method for storing and preserving  
255 tendons in order to minimise changes to the elasticity of the tissue.

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### 323 **Figure Legends**

324 Figure 1 (A) Extracellular matrix of tendons comprising of collagen fibrils embedded in a  
325 hydrated proteoglycan-rich matrix. Decorin core proteins, with glycosaminoglycans  
326 (GAGs) attached, are depicted on the surface of collagen fibrils. (B) A stress-strain curve  
327 (M20) illustrating the areas under the curve corresponding to the strain energy density to  
328 maximum stress ( $u$ ) and to fracture ( $u_p$ ); arrows P1 and P2 indicate the point of inflexion  
329 and the maximum stress ( $\sigma$ ), point respectively.

330 Figure 2. Plot of stress versus strain from control and treated samples. Data points  
331 represent a typical experiment from one sample each.

332 Figure 3. Microscopic examination of tendon samples (a) during initial loading, (b) post-  
333 peak stress,  $\sigma$ . The images in part a and b were taken from a sample belonging to the  
334 M20 (samples treated at  $-20$  °C) group. Horizontal (scale) bar has a length of  $100$   $\mu\text{m}$ .  
335 Arrow A indicates a thin section bridging the bulk of tissue. Arrow B indicates the ends  
336 of the fibres. Arrow C indicates longitudinal splits in the tendon.

337

338 **Tables**

339 Table 1. *Mechanical properties of tail tendons from the C57BL6 mouse*

	$\sigma$ (MPa)	$E$ (MPa)	$u$ (MPa)	$u_p$ (MPa)	$\epsilon$
Control	55.0 $\pm$ 4.9	434.2 $\pm$ 39.5	6.1 $\pm$ 0.7	2.5 $\pm$ 0.9	0.199 $\pm$ 0.008
M20	59.0 $\pm$ 3.4	534.0 $\pm$ 28.3	6.3 $\pm$ 0.5	1.8 $\pm$ 0.2	0.184 $\pm$ 0.005
M80	79.5 $\pm$ 9.7 *	644.5 $\pm$ 71.7 *	10.0 $\pm$ 1.6 *	1.4 $\pm$ 0.2	0.202 $\pm$ 0.011

340 Values are mean  $\pm$  s.e. (n = 10).

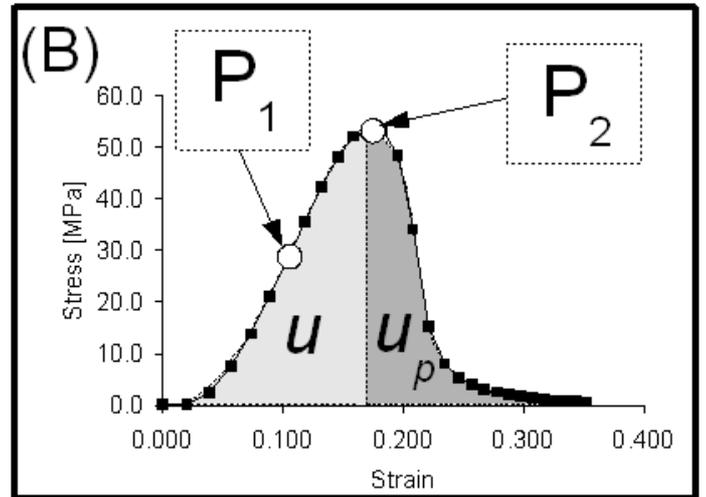
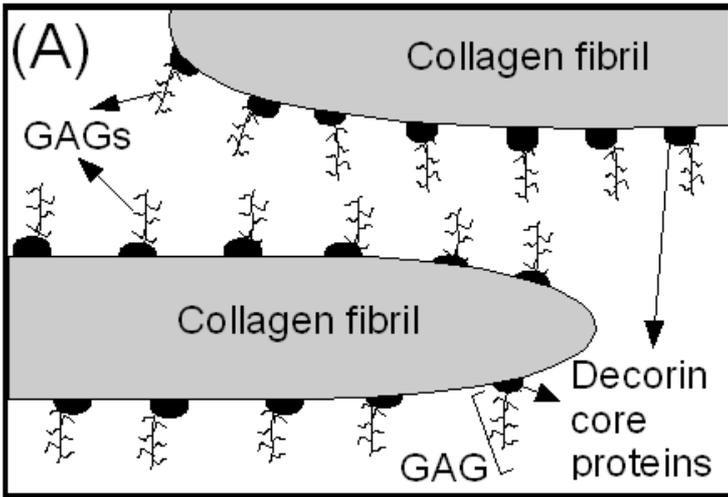
341 \*Significant difference of treatment (with respect to the control) at P < 0.05.

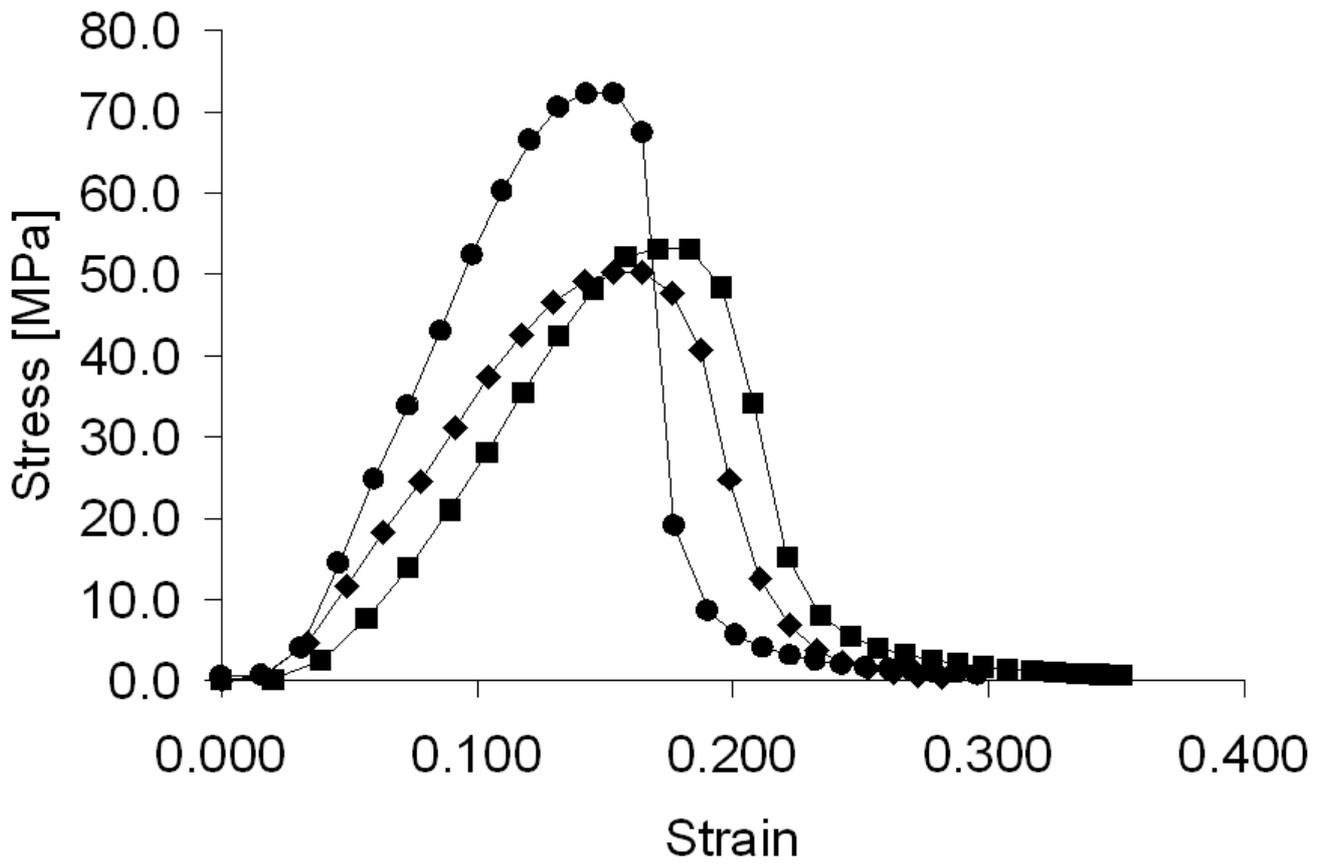
342 Abbreviations are: M20 = samples treated at -20 °C; M80 = samples treated at -80 °C; s.e.

343 = standard error.

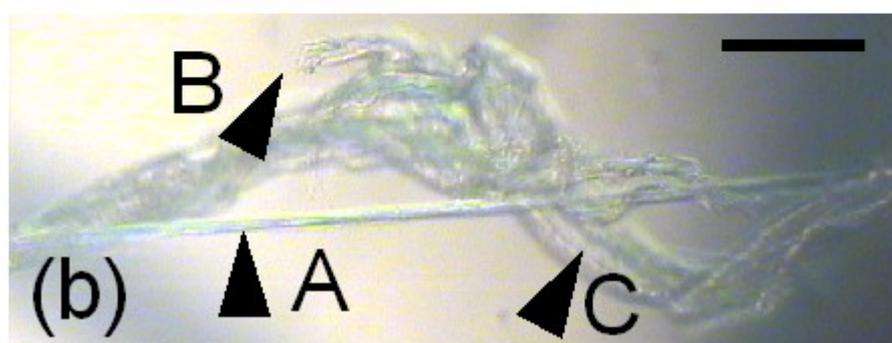
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Legend    —◆— control    —■— M20    —●— M80



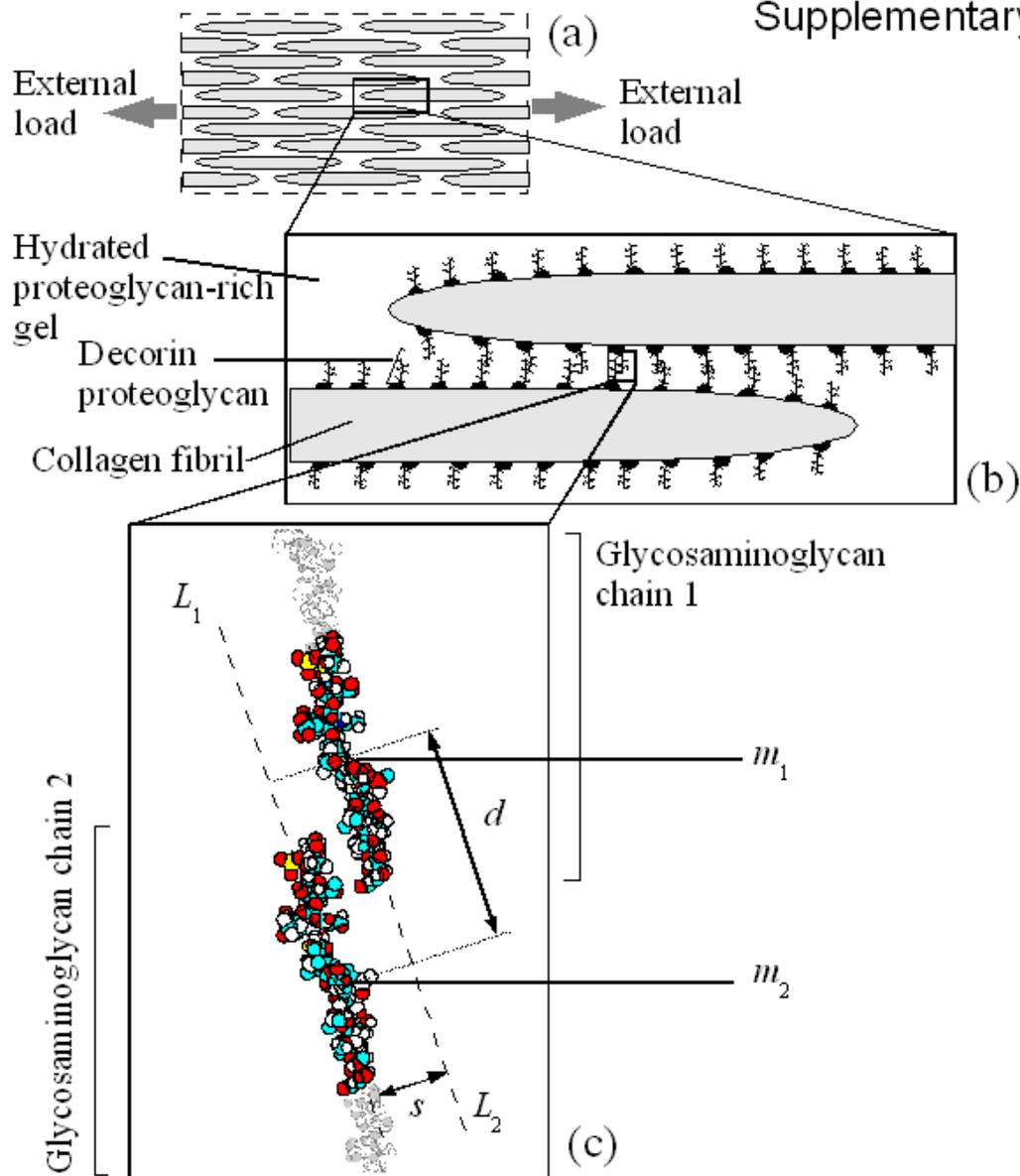


Figure 1S Molecular mechanics study of decorin interactions on adjacent fibrils. A model was developed to evaluate the relationship between VDW energy and cation number density ( $\rho$ ) under fibril-fibril sliding action (a-b). As shown in (c), the model displays the GAG chains (comprising chondroitin-4-sulfate disaccharide repeats) associated with decorin molecules on adjacent fibrils in close proximity. Here,  $s$  and  $d$  denote the separation and displacement between  $m_1$  and  $m_2$  which are the mid-points within the last four repeats of each GAG chain.  $L_1$  and  $L_2$  are the labels for the axes (long dashes) corresponding to GAG chain 1 and 2, respectively.

## Supplementary material

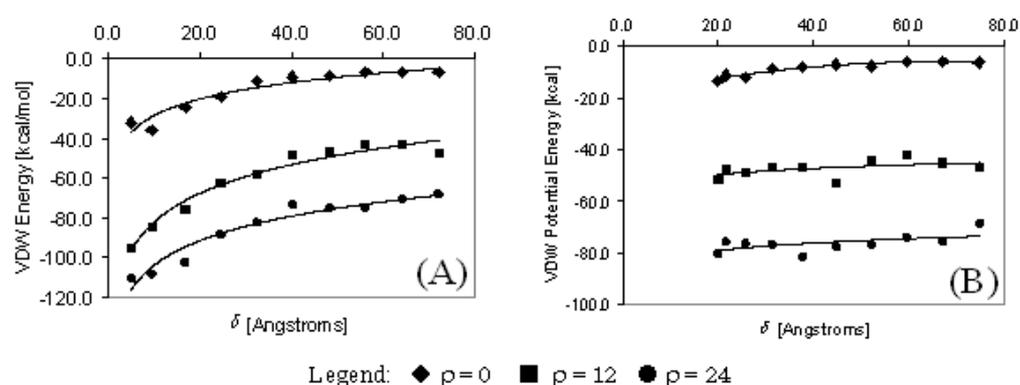


Figure 2S. Graphs of Van Der Waals (VDW) potential energy versus sliding distance  $\delta$  for (A)  $s = 5 \text{ \AA}$  and (B)  $s = 20 \text{ \AA}$ . Here, VDW energy was evaluated from the model (see Figure 1S) as a function of  $\delta$ , where  $\delta^2 = s^2 + d^2$ . An initial condition was used to describe the GAG-GAG overlapping distance over four chondroitin-4-sulfate disaccharide repeats. Each disaccharide repeat is negatively charged due to the anionic nature of the sulphate and carboxylate groups. Restraints on each chain were implemented to maintain the semi-rigid linear conformation of the chain. Solvent environment was modelled implicitly; the dielectric constant was assigned a value of 80 to correspond to water. The model was implemented using software Hyperchem (Version 7.5, Hypercube, Inc.). Geometry optimisation was investigated using molecular mechanics (AMBER3 force field); the Polak Ribiere (conjugate gradient) algorithm was used to determine the minimum energy value. Atomic partial charges were evaluated successively on the sulphate group and carboxyl group, where the total charge value = -1 for each group, using the single point calculation. Mobile cations ( $\text{Na}^+$ ) were designated in closed proximity to each group to model a hypertonic PG matrix environment. The model predicted that VDW energy increases by about 1.2 times as the number of cations increases two-fold. The value of  $\rho$  also influences the rate of decrease in VDW energy with respect to  $\delta$  between the GAGs. The predictions show that VDW energy decreases least rapidly as  $\delta$  increases when  $\rho = 24$ . However, VDW energy decreases most rapidly for  $\rho = 0$ .